

Kinetics of proton pumping in cytochrome *c* oxidase

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We propose a simple model of cytochrome *c* oxidase, including four redox centers and four protonable sites, to study the time evolution of electrostatically coupled electron and proton transfers initiated by the injection of a single electron into the enzyme. We derive a system of master equations for electron and proton state probabilities and show that an efficient pumping of protons across the membrane can be obtained for a reasonable set of parameters. All four experimentally observed kinetic phases appear naturally from our model. We also calculate the dependence of the pumping efficiency on the transmembrane voltage at different temperatures and discuss a possible mechanism of the redox-driven proton translocation. © 2009 American Institute of Physics.

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I. INTRODUCTION

The last enzyme of the respiratory chain of animal cells and bacteria, cytochrome *c* oxidase (CcO), operates as an efficient nanoscale machine converting electron energy into a transmembrane proton electrochemical gradient.^{1–7} The adenosine triphosphate (ATP) synthase enzyme uses this energy to synthesize ATP molecules serving as the “energy currency” of the cell. The process of energy conversion starts when a molecular shuttle, cytochrome *c*, delivers, one by one, high-energy electrons to a dinuclear copper center, Cu_A, located near a positive side (*P*-side) of the inner mitochondrial membrane (see Fig. 1). In recent time-resolved optical and electrometric studies⁸ of the CcO transition from the oxidized (*O*) state to the one electron reduced form (*E*), a *single* electron is donated to the Cu_A redox center by a laser-activated molecule of ruthenium bispyridyl (RubiPy). Thereafter, in a few microseconds ($\sim 10 \mu\text{s}$), a major part of an electron density ($\sim 70\%$) is transferred from the Cu_A center to the low-spin heme *a* (Fe-*a*). Heme *a* is located within the membrane domain at a distance about $2/3$ of the membrane width, W , counting from the *N*-side.^{8–10} Within a time interval of approximately $150 \mu\text{s}$, about 60% of the electron population is transferred from heme *a* to heme *a*₃ (Fe-*a*₃). Heme *a*₃, jointly with the next electron acceptor in line, a copper ion Cu_B, form a binuclear center (BNC) (sites *R* and *B* in Fig. 1), serving as an active catalytic site for dioxygen reduction to water. The redox centers *a*, *a*₃, and Cu_B are located approximately at the same distance ($2/3W$) from the *N*-side of the membrane as heme *a*.

The next phase (with a time scale of the order of $800 \mu\text{s}$) is characterized by a complete electron transfer to the copper ion Cu_B. Time-resolved measurements⁸ show that the first “ $10 \mu\text{s}$ ” phase of the electron transfer process is not

accompanied by a proton transfer, but the “slowness” of the second “ $150 \mu\text{s}$ ” and the third “ $800 \mu\text{s}$ ” phases hints to the proton participation during phases.

The proton path from the negative side of the membrane (*N*-side) toward the *P*-side (for pumped protons) and toward the BNC (for substrate or “chemical” protons) goes through the residue *E278* (for the *Paracoccus denitrificans* enzyme⁸). These residues are located at the end of the so-called *D*-pathway (Fig. 1). A fraction of the substrate protons can also be delivered to the BNC via an additional *K*-pathway, which we will not consider here. The proton to be pumped is supposed to move from *E278* (schematically shown as site *D* in Fig. 1) to an unknown protonable “pumping” site *X* (likely a heme *a*₃ propionate), located above the BNC,¹¹ and, thereafter, via an additional protonable site *C*,^{12,13} to the *P*-side of the membrane. After a fast reprotonation from the *D*-channel, the residue *E278* can donate a substrate proton to the catalytic site near the BNC [probably, to an OH[−] ligand of Cu_B (Refs. 8 and 9)]. It is assumed^{8,9} that during the second $150 \mu\text{s}$ phase, the first (prepumped) proton moves from the residue *E278* (site *D* in Fig. 1) to the pump site *X*, whereas in the third $800 \mu\text{s}$ phase, the second (substrate or chemical) proton populates a catalytic site *Z* near the BNC. In the final phase, which occurs in 2.6 ms, the first proton (in *X*) is translocated (via *C*) to the *P*-side, which is characterized by a higher electrochemical potential than the *N*-side of the membrane.

As a result of all these processes, two protons are taken from the *N*-side of the membrane, and one electron is taken from the *P*-side, and eventually one proton is pumped to the *P*-side. Moreover, one proton and one electron are consumed at the catalytic site to finally produce a water molecule around the BNC. It should be noted that kinetic phases with

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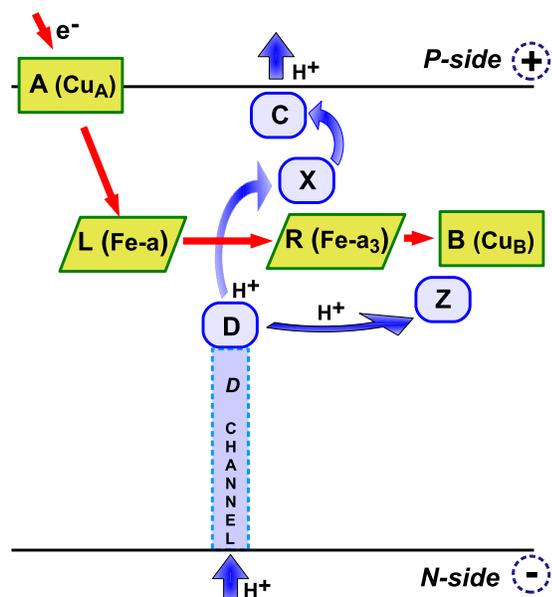


FIG. 1. (Color online) Schematic of CcO. A single electron enters the enzyme at site *A* and travels subsequently to sites *L*, *R*, and, finally, to site *B*. Protons, taken at the *N*-side of the membrane, move to site *D*, which can donate protons both to the pump site *X* and to the catalytic site *Z*. The prepumped proton is transferred from site *X* to the *P*-side of the membrane via site *C*.

similar time scales ($10 \mu\text{s} \rightarrow 100 \mu\text{s} \rightarrow 1000 \mu\text{s}$) have been revealed in other CcO enzymes at various transition steps between the states of the enzyme.^{12,14,15}

Kinetic data obtained in experiments^{8,12,14,15} reflect important details of the still elusive proton pumping mechanism in CcO. To extract these details and gain a deeper insight into the operating principles of the CcO proton pump, it is necessary to compare results of experiments with theoretical predictions. In Ref. 16, a simplified empirical valence bond (EVB) effective potential was combined with a modified Marcus equation to model time-dependent electron and proton transfers in CcO in the range of milliseconds. However, this approach was applied to the single transfer event, not to the sequence of events, and the obtained time scale ($1 \mu\text{s}$) differs by orders of magnitude from the experimental data (about $100 \mu\text{s}$). A computational analysis of the CcO energetics was presented in Refs. 17–21 with molecular models reproducing energetic barriers for the proton transfer steps.^{17,18} The obtained energetic map of the proton and electron pathways in the CcO enzyme can be converted into a set of rate constants, which qualitatively explains the kinetics and unidirectionality of the pumping process. However, these studies do not result in a quantitative model of the efficient CcO proton pump. Moreover, the error range of these semimicroscopic calculations ($\sim 2 \text{ kcal/mol}$) is sometimes higher than the difference between the energy barriers.^{18,20}

Kinetic models of the proton pumping process were also discussed in Ref. 22. Within the master equation approach, it was shown that the proton pumping effect can be achieved in a simplified system having one redox and two proton sites and, with a higher efficiency, $\eta \sim 0.9$, for the design with two redox and two protonable sites, which are electrostatically coupled to each other. However, this work does not contain

any predictions for the kinetics of the pumping process in more realistic setups, with at least four redox sites (Cu_A , heme *a*, heme *a*₃, and Cu_B) and two protonable sites (a residue *E278* and a pump site *X*). To find proper parameters for the proton pump, the authors of Ref. 22 resort to a Monte Carlo search in a multidimensional parameter space. It is hard to imagine, however, that a random search can provide a reasonable set of parameters which will comply with all physical restrictions of real pumps. In general, for comprehensive theoretical studies, it is preferable to determine the relevant parameters of the system using detailed microscopic calculations (see, e.g., Refs. 16–18). However, the huge computational complexity of biological structures makes such an approach extremely difficult. In our paper, we include reasonable estimates for the system parameters into a model describing almost-simultaneous electron and proton transfer processes and compare the obtained kinetics to experimentally observed time scales and site populations of CcO.⁸

The time evolution of the proton pumping process in CcO, related to the experimental data of Refs. 8, 14, 12, and 15, was discussed in Refs. 9, 10, 19, 20, and 23. In these works, the kinetics of the electron-proton system is broken down into a cascade of quasiequilibrium states characterized by distributions of electrons and protons over the sites, as well as by a set of transition rates corresponding to specific kinetic phases. It should be emphasized, however, that many electron and proton transfers can be separated by only a nanosecond time scale, and consequently, the experimentally observed kinetic rates comprise contributions of several almost-simultaneous individual electron and proton transfer events.^{10,23} Correspondingly, an approach taking into account the kinetic inseparability of electron and proton transitions can be useful for understanding recent experimental findings.⁸ We note that kinetic coefficients used in the theoretical analysis of Refs. 10, 19, 20, and 23 were deduced from experiments without independent microscopic calculations of the heights of individual electron and proton barriers.

In the present paper, we analyze electron and proton kinetics in CcO within a simple physical model including four redox centers and four protonable sites electrostatically coupled to each other in the presence of a dissipative environment. Using the master equation approach, we reproduce all four kinetic phases observed in Ref. 8 for a reasonable set of parameters. It should be emphasized that we have performed extensive numerical studies for a wide range of parameters and we found that our model of proton pumping is quite robust to significant variations of the parameters. The specific set of parameters presented below gives a very good agreement with the experimental data of Ref. 8. We consider a single cycle of events, which starts at $t=0$ with one electron transfer to the Cu_A center and finishes at the moment $t=t_B$, when the redox site Cu_B is completely reduced. Notice that the injection of additional high-energy electrons is necessary to maintain this nonequilibrium state of the CcO enzyme. We also determine the efficiency of the proton pumping for our model and its dependencies on the temperature and transmembrane voltage.

The rest of paper is structured as follows. Our model and its parameters are presented in Sec. II. Results of numerical studies are shown in Sec. III and discussed in Sec. IV. Section V contains the conclusions of our work. The detailed derivation of the master equations and the measurable variables is presented in Appendix A. It should be noted that while the results of this paper are obtained in the *classical* regime, our approach (based on quantum transport theory) can be used to examine fine *quantum* effects and consequently, the detailed derivation is worth presenting here.

II. MODEL

As in the real CcO enzyme,^{24–28} the redox chain of the present model includes four centers: Cu_A (site A), heme *a* (Fe-*a*, site L), heme *a*₃ (Fe-*a*₃, site R), and Cu_B (site B), as schematically shown in Fig. 1. The transport chain for protons has four sites: *D* (presumably related to the residue E278 near the end of the *D*-pathway), *X* (the pump site above the BNC), a protonable site *C* placed on the way from the *X*-site to the *P*-side of the membrane, and finally, a protonable site *Z* located in the proximity of the BNC and related to the OH[−] ligand of Cu_B (see Fig. 1). Sites *B* and *Z* serve as final destinations for the injected electron and for the substrate proton, respectively. We assume that the electron can be transferred between the pairs of redox states *A* and *L*, *L* and *R*, *R* and *B*; and that protons can be translocated between the pairs of protonable sites *D* and *X*, *X* and *C*, as well as *D* and *Z*.

To provide an “openness” of the CcO enzyme, which is inherent in the living systems,²² we allow proton transitions between site *D* and the negative side of the membrane as well as between site *C* and the positive side of the membrane. Protons are delivered to the catalytic site *Z* partially through the *K*-pathway.^{5,9} This channel can be incorporated into our model, but for simplicity, it will be neglected. The *N*- and *P*-sides of the membrane play roles of proton reservoirs which work as a source (*N*-side) and a sink (*P*-side) of protons for the enzyme. The redox sites are disconnected from electron reservoirs, and *only one* electron is injected into the redox chain at the initial moment of time, $t=0$.

With the condition of single occupation of each individual site, the system can be populated with up to four protons. Following the setup of Ref. 8, we assume that CcO is populated with a single electron initially located on site A. To quantitatively describe this system we introduce 64 basis states $|\mu\rangle, \mu=1, \dots, 64$ (see Appendix A). The time evolution of the probability distribution over the basis states, $\langle\rho(t)\rangle=\{\langle\rho_\mu(t)\rangle\}$, is governed by the system of master equations, Eq. (A32), with the solution given by Eq. (A34) in the Appendix A. The time-dependent probability distribution $\langle\rho(t)\rangle$ allows us to determine the average populations of all electron and proton sites, $\langle n_\alpha \rangle$ and $\langle N_\beta \rangle$, as functions of time. We can also calculate the number of protons, $\langle N_P(t) \rangle$, translocated to the positive side of the membrane [see Appendix A, Eq. (A36)]. The value of $\langle N_P \rangle$ taken at the end of the pumping cycle ($t=t_B$) determines the pumping efficiency η defined²² as the number of protons pumped across the membrane per electron consumed:

$$\eta = \langle N_P(t_B) \rangle. \quad (1)$$

Note that the efficiency η can take negative values in the case when protons move back from the positive side to the negative side of the membrane.

A. Electrostatic interaction

The electrostatic interaction between the redox ($\alpha = A, L, R, B$) and protonable ($\beta = D, X, C, Z$) sites plays a pivotal role in the electron-proton energy exchange. It should be noted that we consider here only *direct* Coulomb interactions between electron and proton subsystems and between protons themselves. This removes strict geometrical restrictions on the relative positions of electron and proton active sites imposed in our previous model²⁹ based on the Förster-type energy exchange between electrons and protons. Microscopic calculations of the electrostatic parameters, $u_{\alpha\beta}$ and $u_{\beta\beta'}$, involved in the Hamiltonian H_0 [see Appendix A, Eq. (A1)], require a detailed knowledge of the CcO structure complemented by the comprehensive dielectric map of the enzyme.^{17,30,31} Instead, we tune the Coulomb energies to get the best possible fitting of the time scales and site populations measured in the experiment.⁸ The obtained values of Coulomb parameters correlate well with information about the distances between the active sites^{24–28} for reasonable values of the effective dielectric constants.

To describe the experimentally observed kinetic phases of the pumping process, we assume that the coupling, $u_{BZ} = 630$ meV, between the copper ion Cu_B and the catalytic site *Z* [likely an OH[−] ligand of Cu_B (Refs. 8 and 9)] and the coupling, $u_{RX} = 555$ meV, between heme *a*₃ and the pump site *X* are higher than the electrostatic energies $u_{RZ} = 530$ meV and $u_{BX} = u_{XZ} = 510$ meV. Structural studies of the CcO enzyme^{24–28} performed at a resolution of about 2 Å show that the BNC redox sites *R* (heme *a*₃), *B* (Cu_B) and the protonable sites *X* and *Z* are separated by a distance of the order of 6 Å. The value of the electrostatic coupling between these sites, $u \sim 600$ meV, roughly corresponds to the effective dielectric constant, $\epsilon \sim 4$, which is of frequent use for a description of a dry protein interior.^{17,20,21} It should be emphasized, however, that the concept of dielectric constant is not completely appropriate for a calculation of Coulomb potentials in the heterogeneous environment inside and near the BNC.^{30,31}

The distances, r_{LD}, r_{RD} , between the residue E278 (site *D*) and sites *L* and *R* are almost the same: $r_{LD} = 12.3$ Å and $r_{RD} = 12.8$ Å.^{28,32} We estimate the electrostatic coupling between these sites as $u_{LD} \sim u_{RD} \approx 75$ meV, which corresponds to the higher dielectric constant $\epsilon \sim 15$. We consider a smaller dielectric constant, $\epsilon \sim 10$, for the interaction, $u_{LX} = 100$ meV, between sites *L* and *X* separated by the distance $r_{LX} \sim 14.2$ Å.²⁰ Distant-dependent dielectric constants, $\epsilon(r_{\alpha\beta})$, are common in protein electrostatics.^{17,30,31}

Note that here, as in the models of Refs. 9 and 19, the electrostatic coupling, u_{RX} , between heme *a*₃ (site *R*) and site *X* is stronger than the interaction, u_{LX} , between heme *a* (site *L*) and the pump site *X*. For the other parameters we choose the following values (in meV): $u_{DX} \sim 60$, $u_{DZ} \sim u_{BD} \sim 70$, $u_{XC} \sim 100$, $u_{AD} \sim 25$, $u_{AZ} \sim 20$. The Coulomb ener-

gies u_{CZ} , u_{DC} , u_{AX} , u_{AC} , u_{RC} , u_{LC} , and u_{LZ} are assumed to be near 30 meV. Despite the fact that these energies are about or higher than the temperature energy scale, $T=298$ K \sim 26 meV, they have a minor influence on the performance of the model.

B. Energy levels of the sites

We assume that the difference $\Delta\mu$ (A37) between the electrochemical potential μ_P of the P -side and the potential μ_N of the N -side of the membrane is about 210 meV at standard temperature, $T=298$ K, with $\mu_P=105$ meV and $\mu_N=-105$ meV. This corresponds to voltage $V\approx 150$ meV applied across the membrane. We include the electron charge in the parameter V and measure voltage, along with other energies, in units of meV. According to Eq. (A38), the energy levels, ε_α and ε_β , of the electron and proton centers are shifted from their intrinsic values $\varepsilon_\alpha^{(0)}$ and $\varepsilon_\beta^{(0)}$ depending on the voltage V and on the positions x_α, x_β of the active sites. To estimate the electron and proton energies, we take into account the facts⁶ that cytochrome c delivering electrons to the CcO enzyme has a redox potential of order of 250 meV, and that the total drop of electron energy between cytochrome c and the dioxygen reduction site B is about 550 meV. The equilibrium midpoint potentials^{8,9} of the Cu_A center ($E_m\approx 250$ meV) and heme a ($E_m\approx 270$ meV) can also be used as a general guide for estimating energies,³³ although the real parameters can deviate from the estimated values.

We find that our model performs with the high efficiency, $\eta\sim 0.95$, and reproduces all experimentally observed kinetic phases⁸ for the following set of electron intrinsic energies (in meV): $\varepsilon_A^{(0)}=-175$, $\varepsilon_L^{(0)}=-240$, $\varepsilon_R^{(0)}=-185$, $\varepsilon_B^{(0)}=-155$, and for the following energies of protonable sites (in meV): $\varepsilon_D^{(0)}=-100$, $\varepsilon_X^{(0)}=250$, $\varepsilon_C^{(0)}=195$, and $\varepsilon_Z^{(0)}=-65$. It should be noted that in the presence of the transmembrane voltage, $V=150$ meV, the electron energy levels of A and L sites, $\varepsilon_A=-250$, $\varepsilon_L=-265$, are close to the values extracted from equilibrium redox titrations [see also Ref. 34, where an estimation, $(\varepsilon_A-\varepsilon_L)\approx 18$ meV, has been obtained]. For energies of other redox sites we use the values: $\varepsilon_R=-210$, $\varepsilon_B=-180$. The energies of the protonable sites are also shifted with voltage, $V=150$ meV, present: $\varepsilon_D=-85$, $\varepsilon_X=295$, $\varepsilon_C=270$, and $\varepsilon_Z=-40$. It should be stressed that the energy, ε_X , of the pump site X is set to be higher than the potentials of the proton reservoirs on both sides of the membrane: $\varepsilon_X>\mu_P>\mu_N$. However, the presence of an electron on site R decreases the proton energy to the level, $\varepsilon\sim(\varepsilon_X-u_{RX})\sim -260$ meV, which is below the energy of the D -site and below the electrochemical potential, $\mu_N=-105$ meV, of the N -side of the membrane. As a result, the pump site X is populated with a prepumped proton. When the chemical proton moves to site Z and the electron is transferred to the B -site, the energy level of the X -site returns to the initial position, $\varepsilon_X=295$ meV, since the electron and proton charges of the catalytic site compensate each other, $u_{BX}=u_{XZ}$. The high-energy prepumped proton can now move to site C and after that, to the P -side of the membrane characterized by the potential $\mu_P=105$ meV. A large energy gap,

$(\varepsilon_X-\varepsilon_D)\sim 380$ meV, significantly suppresses the return of the X -proton to site D and to the N -side of the membrane.

C. Reorganization energies and transition rates

Part of the energy delivered to the redox center Cu_A at the initial time, $t=0$, is dissipated to an environment characterized by sets of electron ($\lambda_{\alpha\alpha'}$) and proton ($\lambda_{\beta\beta'}$) reorganization energies. To be efficient, the proton pumping process should occur with minimal energy dissipation. It is shown in Ref. 35 that the reorganization energy for the a to a_3 electron transfer in the CcO enzyme can be as low as 100 meV. Similar estimates apply for the proton reorganization energies.^{36,37} Here, we use the higher energy parameter, $\lambda_{AL}=200$ meV, for the A -to- L transfer and accept the lower value, $\lambda_{\alpha\alpha'}\approx\lambda_{\beta\beta'}\approx 100$ meV, for other electron and proton transitions. It is argued in Refs. 34, 38, and 39 that for the $\text{Cu}_A\rightarrow$ heme a electron transition the reorganization energy must be in the range from 150 to 500 meV, which is much lower than the typical values of the reorganization energy for electron transfers in protein. The low values of electron reorganization energies ($\lambda\sim 2-4$ kcal/mol) have also been calculated for electron transfer reactions in *Rhodobacter sphaeroides*.³⁰

To reproduce the initial kinetic phases, we use the following tunneling energies: $\Delta_{AL}\sim 0.9$ μeV and $\Delta_{LR}\sim\Delta_{BR}\sim 14$ μeV . The parameters Δ_{LR} and Δ_{BR} describe the electron transfers, which are coupled to the slower proton transitions characterized by the energy scales: $\Delta_{DX}\sim\Delta_{CX}\sim 0.3$ μeV and $\Delta_{ZD}\sim 0.2$ μeV . It should be noted that the electron transfer between hemes a and a_3 can occur in a nanosecond time scale.⁴⁰ The hydrogen-bonded chains in proteins are also able to conduct protons in nanoseconds or faster.^{41,42}

We also select the values $\Gamma_N\sim\Gamma_P\sim 17$ ms^{-1} for the parameters Γ_N and Γ_P , which determine the flow of protons through the enzyme. These parameters Γ_N and Γ_P are of the same order as some of the transition rates $\kappa_{\mu\nu}$ used in Ref. 22.

III. RESULTS

A. Four kinetic phases

In Fig. 2, starting at $t=0.1$ μs , we show a process of population and depopulation of the electron, $n_{A,L,R,B}$, and proton, $N_{D,X,C,Z}$, sites as well as the time dependence of the average number of protons pumped to the positive side of the membrane, N_p . From here on we drop the brackets $\langle\dots\rangle$ denoting the averaging over the environmental fluctuations and over the states of the proton reservoirs. The calculations are performed for the standard conditions ($\mu_P=105$ meV, $\mu_N=-105$ meV, $\Delta pH=-1$, and $T=298$ K) and for the transmembrane voltage $V=150$ meV. We assume that initially a single electron is located at site A (Cu_A), and a proton occupies site D . This means that at $t=0$ only one element of the density matrix is not equal to zero: $\rho_2(0)=1$.

During the first phase of the process the electron moves from site A to site L (heme a). In ~ 10 μs near 70% of the electron density is transferred to heme a (site L) with the remaining 30% distributed almost equally between site A

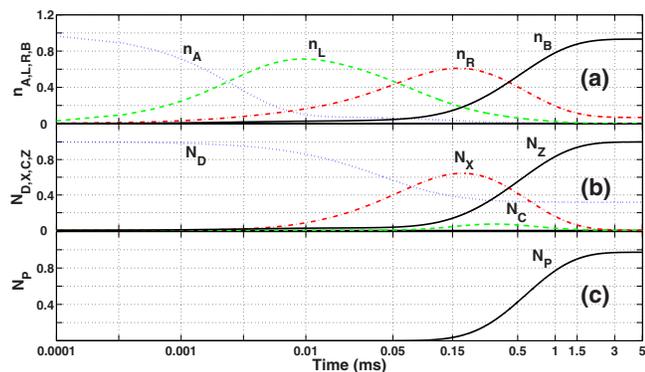


FIG. 2. (Color online) Kinetics of electron and proton transfers in CcO for $\mu_p=105$ meV, $\mu_N=-105$ meV, $T=298$ K, and $V=150$ meV. The time axis (in ms) is shown in a logarithmic scale starting at $t=0.1$ μ s. The process begins at $t=0$, when a single electron populates site A, and a single proton is located on site D. (a) Time dependence of the electron populations n_A (blue dotted line), n_L (green dashed line), n_R (red dashed-dotted line), and n_B (black continuous line). (b) The proton populations, N_D (blue dotted line), N_X (red dashed-dotted line), N_C (green dashed line), and N_Z (black continuous line), vs time. (c) The number of pumped protons, N_p , as a function of time. The first phase of the process, where the electron moves from site A to site L, corresponds to the maximum of the population n_L at the moment $t \approx 10$ μ s. In the second phase, both the electron population of the R-site, n_R , and the proton population of the X-site, N_X , peak at $t \approx 150$ μ s. The third phase of the process is marked by the significant population of the proton site Z and the electron site B at $t \sim 1$ ms. In this phase site X is depopulated, and the prepumped proton is partially transferred to the P-side of the membrane, $N_p \sim 0.8$. In the final phase, at $t \sim 3$ ms, the electron site B and the catalytic proton site Z are occupied, $n_B \approx 1$, $N_Z = 1$, and about one proton ($N_p \sim 0.95$) is translocated to the P-side of the membrane.

(Cu_A) and site R (heme a_3). This corresponds roughly to the 70% electron population of heme a after the first 10 μ s phase observed experimentally in Ref. 8. No pronounced changes in populations of the protonable sites accompany this stage [see Fig. 2(b)].

The second phase of the electron transfer is postponed by the time 150 μ s, despite the fast intrinsic transition rate between the L and R redox sites. Besides the 55 meV potential difference between sites R and L, the electron transfer in this phase is hampered by the involvement of protons. It is evident from Figs. 2(a) and 2(b) that with a microsecond delay, the slightly uphill electron transfer from site L to site R is followed by the proton translocation from site D ($E_D = -85$ meV) to the pump site X having much higher initial energy, $E_X = 295$ meV. This transition has been made possible by the strong R-X Coulomb attraction ($u_{RX} = 555$ meV) lowering the effective energies of both electron and proton sites. In line with the experimental data⁸ at $t = 150$ μ s, the electron density is located mainly on site R (60%) and partially on site L ($\sim 20\%$), and on site B ($\sim 15\%$). Site A is practically empty at this stage. It is important that at almost the same moment of time ($t \approx 150$ μ s) the population of the protonable pump site X also reaches its maximum ($\sim 65\%$).

It is evident from Fig. 2(b) that the occupation of the pump site X is accompanied by the monotonic population of the protonable catalytic site Z, thus lowering the energy of the redox site B from its initial level, $\epsilon_B = -180$ meV, to the final value of the order of -820 meV (see also Fig. 3). The population of B-site, n_B , closely follows (with a small delay)

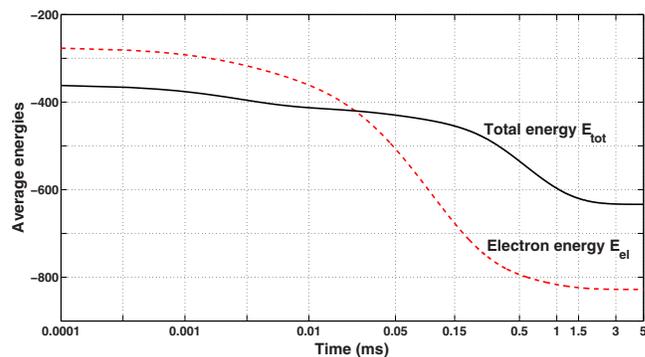


FIG. 3. (Color online) The total energy of the system, E_{tot} (black continuous line), and the average energy of the electron, E_{el} , as functions of time (in ms, logarithmic scale). The electron energy is varied in the range from $E_{el} = -277$ meV at $t \approx 0$, to the value $E_{el} = -828$ meV at $t = t_B = 5$ ms. The whole electron-proton system dissipates less energy, $\Delta E_{tot} \approx 270$ meV, than its electron component, which loses about 550 meV during the pumping process, indicating the energy transfer to the proton subsystem.

the population N_Z of the proton catalytic site Z [see Figs. 2(a) and 2(b)]. It can be seen from Fig. 2(b) that in ~ 300 μ s the pumped proton moves from site X to the transient site C, placed between X and the P-side of the membrane, and after that to the positive side of the membrane.

In the third phase ($t \sim 1$ ms), the substrate (chemical) proton [Fig. 2(b)] occupies the catalytic site Z, $N_Z > 0.8$. Then, with a microsecond delay, the electron [Fig. 2(a)] is transferred, $n_B \geq 0.8$, to the B-center (Cu_B), so that heme a is practically reoxidized, $n_L \sim 0.02$. This stage is correlated with the 800 μ s phase mentioned in Ref. 8.

In the fourth phase ($t \sim 3$ ms), the pumped proton [Fig. 2(c)] moves to the positive site of the membrane, $N_p \approx 0.95$, the substrate proton populates site Z, $N_Z = 1$, and the electron is almost completely transferred to site B, $n_B \approx 1$. On average, about 1.3 protons are taken from the N-side of the membrane during the whole process.

The variations of the average electron energy, $E_{el} = \langle H_{el} \rangle$, and the total energy of the system,

$$E_{tot} = \langle H_0 \rangle + \mu_p N_p + \mu_N N_N,$$

with time are shown in Fig. 3. Here H_0 is the basic Hamiltonian of system (A1), H_{el} is the Hamiltonian of the electron component (A2), N_p and N_N are the average numbers of protons [Eq. (A36)] translocated to the P- or N-side of the membrane, respectively. At the beginning, the electron has energy

$$E_{el}(0) \approx (\epsilon_A - u_{AD}) \approx -277 \text{ meV},$$

and at the end of the process its energy sinks to the level

$$E_{el}(5 \text{ ms}) \approx (\epsilon_B - u_{BZ} - u_{BD}) \approx -828 \text{ meV}$$

with the total drop $\Delta E_{el} \approx 551$ meV, corresponding to the experimental value.⁶ The total energy of the system, E_{tot} , shows a decrease in the order of $\Delta E_{tot} \approx 271$ meV, which is less than the drop of electron energy since one proton gains the energy during its pumping to the positive side of the membrane.

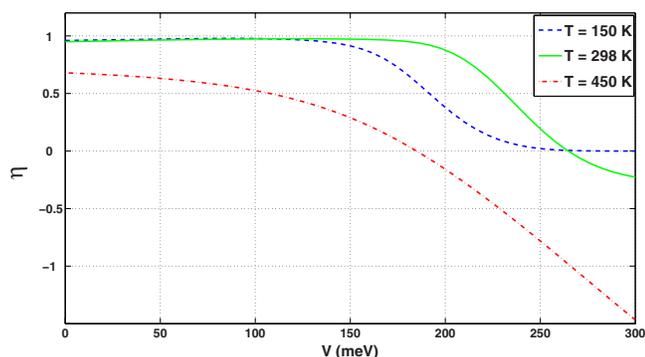


FIG. 4. (Color online) The efficiency of the pump, $\eta = N_p(t_B)$, where $t_B = 5$ ms, as a function of the transmembrane voltage V at the temperatures $T = 150$ K (blue dashed line); $T = 298$ K (green continuous line), and $T = 450$ K (red dashed-dotted line). For the physiological range of transmembrane voltages, $150 \text{ meV} < V < 200 \text{ meV}$, the pump operates with a maximum efficiency at room temperatures ($T = 298$ K). At high temperatures and high enough transmembrane voltages, the efficiency η takes negative values, suggesting that at these conditions the protons flow back, from the positive to the negative side of the membrane.

B. Pumping efficiency

It follows from Fig. 2(c) that at $t = t_B = 5$ ms, the average number of pumped protons, N_p , reaches its peak value, which can be used as a definition²² of the pumping efficiency η : $\eta = N_p(t_B)$. According to this definition, the present model demonstrates an almost-perfect performance with an efficiency $\eta \approx 0.95$ at $T = 298$ K, $\Delta\mu = 210$ meV, and $V = 150$ meV. This is comparable to the efficiency of CcO (Refs. 3 and 6) pumping one proton across the membrane per one electron consumed at the oxygen reduction site. We find that the definition of the efficiency η introduced above is not sensitive to the choice of the specific moment $t_B = 5$ ms since the number of pumped protons, $N_p(t)$, does not decrease noticeably with time during the interval from 5 ms to more than 100 ms at the standard conditions.

In Fig. 4 we plot the pumping efficiency η versus the transmembrane voltage V at three different temperatures: $T = 150$ K (blue dashed line), $T = 298$ K (green continuous curve), and $T = 450$ K (red dashed-dotted line). We assume that the electrochemical gradient $\Delta\mu$ varies in accordance to Eq. (A37) where $\Delta\rho H = -1$. At $T = 150$ K the pumping efficiency η is almost constant at low voltages, $V < 150$ meV, with a subsequent drop at high voltages. The pump works better at room temperatures, $T = 298$ K, and keeps the efficiency steady up to voltages $V \sim 200$ meV. Notice that in this case the efficiency η , which is proportional to the average number of pumped protons, becomes negative at $V \geq 270$ meV. The performance of the model is significantly deteriorated at high temperatures, $T = 450$ K, when the proton flow is reversed starting with the relatively low voltage gradient, $V \sim 180$ meV.

IV. DISCUSSION

The obtained time evolution of the electron and proton populations (see Fig. 2) features four experimentally observed phases of the proton pumping process: The first $10 \mu\text{s}$ phase, when the electron is transferred from Cu_A (site A) to heme a (site L); the second $150 \mu\text{s}$ phase when the

electron moves from heme a to heme a_3 (site R), and, with a microsecond delay, a proton partially occupies the pump site X; the third “ $1000 \mu\text{s}$ ” phase when the chemical proton is transferred to the catalytic sites Z and, a slightly later, the electron is transferred to the ultimate electron acceptor Cu_B . In the fourth phase, at $t \sim 3$ ms, the prepumped proton is released to the P-side of the membrane.

It should be emphasized that contrary to the models proposed in Refs. 8, 9, and 23, this process cannot be described as a sequence of transitions between clearly defined quasi-equilibrium states since many electron and proton transfers occur in a very short time one after the other. The present theoretical model, which includes four redox sites (two copper centers and two hemes) and four protonable sites, is able to explain the efficient performance ($\eta \sim 0.95$) of the real CcO (Ref. 3) pumping almost one proton per one electron consumed against the electric potential difference, $V \geq 150$ meV, and against the transmembrane electrochemical gradient, $\Delta\mu \geq 210$ meV. We stress that all four kinetic phases appear naturally in our model for a reasonable set of the system parameters without artificial inclusions of consequent transfer processes.

The mechanism of the proton pumping analyzed above is based on the direct electrostatic interaction between the redox and protonable sites, especially between the electron located on site R (heme a_3) and the proton located on the pump site X. The Coulomb coupling between the redox site Cu_B and the protonable catalytic site Z plays a very important role as well. The proton to be pumped is sequentially translocated to the P-side of the membrane from sites X and C. At the beginning of the process these sites are empty since their energy levels are assumed to be higher than the energy levels of the proton source (μ_N and E_D) and the proton drain (μ_P). After the first $10 \mu\text{s}$ phase the energy level of the L-site is slightly (~ 55 meV) lower than the energy level of the R-site. However, an interaction with the environment facilitates the slow electron transfer to site R. The population of site R with the electron is accompanied by the lowering of the X-site energy level followed by the proton translocation from site D to the pump site X. Because of the strong X-R electrostatic attraction, the effective energy of the R-electron drops below the energy of the L-site, which results in the second $150 \mu\text{s}$ phase where the major part ($\sim 60\%$) of the electron density is concentrated on site R, and the pump site X is partially ($\sim 65\%$) populated with a proton. The electron transfer to site R also leads to lowering the energy of the Z-site, thus inducing a monotonous population of the catalytic protonable site Z. No switch redirecting protons to site D or to site X (as proposed in Ref. 43) is needed here because both of these sites can be populated from site D.

It should also be emphasized that these three processes—the electron transfer to the R-site, the occupation of the pump site X, and the translocation of a proton to the Z-site—are strongly correlated in time. The proton transfer to the Z-site digs a deep potential well for the electron at site B, and in the third (~ 1 ms) phase the electron falls into this well. Afterward, the Coulomb attraction between the prepumped X-proton and the electron is almost compensated by the electrostatic repulsion between X and Z protons, and

the energy level of the X -proton returns to its original high value. The reverse translocation of the X -proton to site D is strictly suppressed since now the energy difference between sites X and D ($E_X - E_D \sim 380$ meV) significantly exceeds the reorganization energy λ_{DX} as well as the temperature broadening, $2\sqrt{\lambda_{DX}T}$, of the transition rates in Eq. (A27). However, the prepumped proton can easily move to the slightly (~ 25 meV) lower energy level E_C , and, after this, to the positive side of the membrane characterized by the even lower electrochemical potential $\mu_P = 105$ meV. Our model does not require any nonlinear gates⁴³ to prevent a proton leakage from the positive to the negative side of the membrane.

V. CONCLUSION

We have analyzed a simple model describing the kinetics of the proton pumping process in CcO initiated by a single-electron injection. Within our model, this electron is subsequently transferred along four sites electrostatically coupled to four protonable sites. We have shown that the energy loss by this electron facilitates the proton transfer against the transmembrane voltage from the negative to the positive sides of the membrane with the efficiency $\eta \sim 0.95$. In contrast to previous studies, we have not broken the electron and proton transfers into a series of transitions between the independent quasiequilibrium states but examined inseparable dynamics of the pumping process. We have derived the master equations of motion and solved them numerically for a reasonable set of the system parameters. The obtained time evolution naturally encompasses *all four* experimentally observed kinetic phases.

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APPENDIX A: MASTER EQUATIONS

The kinetics of charge transfer in the CcO enzyme can be described by a set of master equations.^{22,44,45} For completeness we present here a derivation of these equations. We start from the formalism of second quantization,^{29,46,47} even though in this paper we only discuss the classical results, with an examination of quantum coherent effects to be performed in the future. Electrons, located on the redox sites α ($\alpha = A, L, R, B$), are described by the creation and annihilation Fermi operators a_α^+, a_α , and protons located on the protonable sites β ($\beta = D, X, C, Z$) are described by the creation and annihilation Fermi operators b_β^+, b_β . The spin degrees of freedom are neglected; thus, each site can only be occupied by a single particle. The electron population of the α -site, n_α , is expressed as $n_\alpha = a_\alpha^+ a_\alpha$, and for a proton population N_β on site β , we have the relation $N_\beta = b_\beta^+ b_\beta$. Protons on the negative (N) and on the positive (P) sides of the membrane (σ

$= N, P$) are continuously distributed over the space of an additional ‘‘quasiwavenumber’’ parameter q and characterized by the creation and annihilation Fermi operators $d_{q\sigma}^+, d_{q\sigma}$ with the density operator $N_{q\sigma} = d_{q\sigma}^+ d_{q\sigma}$.

1. Hamiltonian of the system

The total Hamiltonian H of the electron-proton system incorporates a basic term,

$$H_0 = \sum_{\alpha} \varepsilon_{\alpha} n_{\alpha} + \sum_{\beta} \varepsilon_{\beta} N_{\beta} + \sum_{\beta\beta'} u_{\beta\beta'} N_{\beta} N_{\beta'} - \sum_{\alpha\beta} u_{\alpha\beta} n_{\alpha} N_{\beta}, \quad (\text{A1})$$

where the first and second terms describe the electron (α) and proton (β) sites with energies ε_{α} and ε_{β} , respectively, and the third and fourth terms are responsible for the Coulomb interaction of protons with each other and the electron, respectively. It should be noted that in our single-electron model there is no interelectron Coulomb interaction. We will also calculate the energy of the electron component, which is determined by the Hamiltonian

$$H_{\text{el}} = \sum_{\alpha} \varepsilon_{\alpha} n_{\alpha} - \sum_{\alpha\beta} u_{\alpha\beta} n_{\alpha} N_{\beta}. \quad (\text{A2})$$

For protons in the N -side and P -side reservoirs we introduce the Hamiltonian

$$H_{NP} = \sum_{q\sigma} \varepsilon_{q\sigma} N_{q\sigma}, \quad (\text{A3})$$

with the energy spectrum $\varepsilon_{q\sigma}$, whereas proton transitions between site D and the N -side of the membrane, and site C and the P -side are given by the transfer Hamiltonian

$$H_{\text{tr}} = - \sum T_{qN} d_{qN}^+ b_D - \sum T_{qP} d_{qP}^+ b_C + \text{H.c.}, \quad (\text{A4})$$

characterized by the coefficients T_{qN} and T_{qP} . The component

$$H_{\text{tun}} = - \sum_{\alpha\alpha'} \Delta_{\alpha\alpha'} a_{\alpha}^+ a_{\alpha'} - \sum_{\beta\beta'} \Delta_{\beta\beta'} b_{\beta}^+ b_{\beta'}, \quad (\text{A5})$$

is responsible for electron tunneling between the pairs ($\alpha\alpha'$) of sites $A-L$, $L-R$, $R-B$, and for proton transitions between the pairs ($\beta\beta'$) of sites $D-X$, $X-C$, and $D-Z$, with the corresponding amplitudes $\Delta_{\alpha\alpha'}$ (for electrons) and $\Delta_{\beta\beta'}$ (for protons), where $\Delta_{\alpha\alpha'}^+ = \Delta_{\alpha'\alpha}$ and $\Delta_{\beta\beta'}^+ = \Delta_{\beta'\beta}$.

Protons are delivered from a solution to the active site D by the water-filled D -channel.^{5,7} It was argued^{13,41,48} that the D -channel contains hydrogen-bonded chains of water molecules, which can convey protons via the Grotthuss mechanism. In this case, the proton transfer is considered as a collective motion of a positive charge through the chain, but not as a motion of an individual proton. According to another point of view (see Refs. 17 and 49–51), the proper orientation of water molecules required for the Grotthuss mechanism is characterized by a much smaller energetic penalty than the electrostatic barriers associated with a proton transfer through the channel. Thus, the dominant contribution to the kinetic rate of the proton transport in proteins is provided by the electrostatic energy. In the present work we model

proton transitions (between site D and the N -side as well as between site C and the P -side of the membrane) by the Hamiltonian H_{tr} (A4) with matrix elements that do not specify the transfer origin.

The transport of protons between the active sites D - X , X - C , and D - Z are described by phenomenological coefficients $\Delta_{\beta\beta'}$ in Hamiltonian (A5). To obtain kinetic rates for proton transitions between active sites in the presence of an environment, we resort to the Marcus formulation of the problem. The relevant approach based on the EVB method⁵² has been developed in Ref. 51. As shown in Refs. 16, 17, and 49, the modified Marcus relations can be successfully applied for modeling the proton transfer steps in CcO.

2. Environment

To take into account the interaction of the electron-proton system with its environment, we introduce a term H_{env} ,

$$H_{env} = \sum_j \frac{p_j^2}{2m_j} + \frac{1}{2} \sum_j m_j \omega_j^2 \times \left(x_j - \sum_{\alpha} x_{j\alpha} n_{\alpha} - \sum_{\beta} x_{j\beta} N_{\beta} - \sum_{\sigma} x_{j\sigma} N_{\sigma} \right)^2, \quad (\text{A6})$$

where $N_{\sigma} = \sum_q N_{q\sigma}$ is the total number of protons in the σ -reservoir ($\sigma = N, P$). The environment is represented as a set of harmonic oscillators^{53,54} with coordinates x_j , momenta p_j , masses m_j , and frequencies ω_j . The shifts $x_{j\alpha}$, $x_{j\beta}$, and $x_{j\sigma}$ define coupling strengths of electrons and protons to the environment. The total Hamiltonian H is the sum of all above-mentioned components,

$$H = H_0 + H_{NP} + H_{tr} + H_{tun} + H_{env}. \quad (\text{A7})$$

With the unitary transformation,

$$U = \exp \left\{ -i \sum_j p_j \left(\sum_{\alpha} x_{j\alpha} n_{\alpha} + \sum_{\beta} x_{j\beta} N_{\beta} + \sum_{\sigma} x_{j\sigma} N_{\sigma} \right) \right\} \quad (\text{A8})$$

the total Hamiltonian, $H' = U^+ H U$, can be transformed to the form

$$H' = H_0 + \sum_{q\sigma} \varepsilon_{q\sigma} N_{q\sigma} + \sum_j \left(\frac{p_j^2}{2m_j} + \frac{m_j \omega_j^2 x_j^2}{2} \right) - \sum_{\alpha \neq \alpha'} Q_{\alpha\alpha'} a_{\alpha}^+ a_{\alpha'} - \sum_{\beta \neq \beta'} Q_{\beta\beta'} b_{\beta}^+ b_{\beta'} - \sum T_{qN} d_{qN}^+ b_D - \sum T_{qP} d_{qP}^+ b_C - \sum T_{qP}^* b_C^+ d_{qP}, \quad (\text{A9})$$

where the operators,

$$Q_{\alpha\alpha'} = Q_{\alpha'\alpha}^+ = \Delta_{\alpha\alpha'} \exp \left\{ i \sum_j p_j (x_{j\alpha} - x_{j\alpha'}) \right\}, \quad (\text{A10})$$

$$Q_{\beta\beta'} = Q_{\beta'\beta}^+ = \Delta_{\beta\beta'} \exp \left\{ i \sum_j p_j (x_{j\beta} - x_{j\beta'}) \right\},$$

describe the effect of the environment on the electron and proton transitions. The protonable site C is located near the P -side of the membrane, and site D is tightly coupled to the

N -side by the D -channel. It is reasonable to assume, therefore, that C -to- P and N -to- D proton transitions have a negligible effect on the equilibrium position of the j -oscillator of the environment: $x_{jC} = x_{jP}$, $x_{jD} = x_{jN}$ so that the corresponding phase factors in Eq. (A10) related to the Hamiltonian $H'_{tun} = U^+ H_{tun} U$ and to the total Hamiltonian (A9) can be omitted.

3. Basis states and eigenenergies

To quantitatively analyze the system with a single electron and with up to four protons we introduce a basis of 64 eigenstates of the Hamiltonian H_0 : $|1\rangle = a_A^+ |0\rangle$, $|2\rangle = a_A^+ b_D^+ |0\rangle$, $|3\rangle = a_A^+ b_X^+ |0\rangle$, $|4\rangle = a_A^+ b_C^+ |0\rangle$, $|5\rangle = a_A^+ b_Z^+ |0\rangle$, $|6\rangle = a_A^+ b_D^+ b_X^+ |0\rangle, \dots, |64\rangle = a_B^+ b_D^+ b_X^+ b_C^+ b_Z^+ |0\rangle$. Here $|0\rangle$ is the vacuum state of the system with no electrons and no protons, $|1\rangle$ is the state with an electron on site A , $|2\rangle$ is the state with an electron on site A and a proton on site D , $|3\rangle$ has one electron on site A and a proton on site X , $|4\rangle$ describes the state with an electron on site A and a proton on site C , and so on. Finally, $|64\rangle$ is the state with a single electron on site B and with one proton on each site D, X, C , and Z (i.e., a total of four protons). The state $|1\rangle$ has the eigenenergy $E_1 = \varepsilon_A$, state $|2\rangle$ has the energy $E_2 = \varepsilon_A + E_D - u_{AD}$, and the last state $|64\rangle$, fully loaded with four protons, has the energy

$$E_{64} = \varepsilon_B + \sum_{\beta=D}^{\beta=Z} (\varepsilon_{\beta} - u_{B\beta}) + u_{DX} + u_{DC} + u_{DZ} + u_{XC} + u_{XZ} + u_{CZ}.$$

The Hamiltonian H_0 (A1) is diagonal in the new basis,

$$H_0 = \sum_{\mu=1}^{64} E_{\mu} |\mu\rangle \langle \mu|. \quad (\text{A11})$$

Other operators may have a nondiagonal form in the new basis, for example,

$$a_{\alpha}^+ a_{\alpha'} = \sum_{\mu\nu} (a_{\alpha}^+ a_{\alpha'})_{\mu\nu} \rho_{\mu\nu},$$

$$b_{\beta}^+ b_{\beta'} = \sum_{\mu\nu} (b_{\beta}^+ b_{\beta'})_{\mu\nu} \rho_{\mu\nu}, \quad (\text{A12})$$

$$b_{\beta} = \sum_{\mu\nu} b_{\beta, \mu\nu} \rho_{\mu\nu},$$

where

$$\rho_{\mu\nu} = |\mu\rangle \langle \nu|. \quad (\text{A13})$$

Here indices μ and ν sweep all integers from 1 to 64.

4. Electron and proton transitions

In addition to the diagonal parts H_0 and H_{NP} , the total Hamiltonian of system H contains the term responsible for the proton transitions between the N -side of the membrane and site D , and between the P -side and site C ,

$$H_{tr} = - \sum (T_{qN} b_{D,\mu\nu} d_{qN}^+ + T_{qP} b_{C,\mu\nu} d_{qP}^+) \rho_{\mu\nu} + \text{H.c.}, \quad (\text{A14})$$

as well as the off-diagonal term H_{tun} describing the tunneling of electrons and the transfer of protons between the active sites,

$$H_{\text{tun}} = - \sum_{\mu\nu} \mathcal{A}_{\mu\nu} \rho_{\mu\nu} - \sum_{\mu\nu} \rho_{\nu\mu} \mathcal{A}_{\mu\nu}^+ \quad (\text{A15})$$

Here the operator $\mathcal{A}_{\mu\nu}$ is represented by a linear combination of the bath operators Q_{AL}, \dots, Q_{ZD} [see Eq. (A10)], multiplied by the nondiagonal ($\mu \neq \nu$) transition matrix elements $(a_A^+ a_L)_{\mu\nu}, \dots, (b_Z^+ b_D)_{\mu\nu}$,

$$\begin{aligned} \mathcal{A}_{\mu\nu} = & Q_{AL} (a_A^+ a_L)_{\mu\nu} + Q_{LR} (a_L^+ a_R)_{\mu\nu} + Q_{RB} (a_R^+ a_B)_{\mu\nu} \\ & + Q_{DX} (b_D^+ b_X)_{\mu\nu} + Q_{XC} (b_X^+ b_C)_{\mu\nu} + Q_{DZ} (b_D^+ b_Z)_{\mu\nu}. \end{aligned} \quad (\text{A16})$$

It should also be noted that operators of the N and P proton reservoirs, d_{qN} and d_{qP} , cannot be completely expressed in terms of the basis operators $\rho_{\mu\nu}$.

5. Derivation of the master equations

A probability $\langle \rho_\mu \rangle$ to find the electron-proton system in the state $|\mu\rangle$ is determined by the diagonal operator $\rho_\mu = |\mu\rangle\langle\mu|$ averaged over the states of the environment and over the distributions of protons on both, N and P , sides of the membrane. The time evolution of the operator ρ_μ is governed by the Heisenberg equation

$$\begin{aligned} i\dot{\rho}_\mu = & [\rho_\mu, H_{tr}] - \sum_\nu \{ \mathcal{A}_{\mu\nu} \rho_{\mu\nu} - \mathcal{A}_{\nu\mu} \rho_{\nu\mu} \} \\ & + \sum_\nu \{ \mathcal{A}_{\mu\nu} \rho_{\mu\nu} - \mathcal{A}_{\nu\mu} \rho_{\nu\mu} \}^+. \end{aligned} \quad (\text{A17})$$

To derive a master equation for the probabilities $\langle \rho_\mu \rangle$, we have to average Eq. (A17) and calculate the correlation functions $\langle \mathcal{A}_{\mu\nu} \rho_{\mu\nu} \rangle$ of the environment operators $\mathcal{A}_{\mu\nu}$ (A16) and the operators of the system $\rho_{\mu\nu}$. The transition coefficients, $\Delta_{\alpha\alpha'}$, $\Delta_{\beta\beta'}$ and $T_{q\sigma}$, are supposed to be much smaller than the energy scales given by the basis spectrum E_μ ($\mu = 1, \dots, 64$). This means that the effective interactions with the N and P proton reservoirs [see Eq. (A14)] and with the bath of oscillators [see Eqs. (A15) and (A16)] can be treated as a perturbation. In the framework of the theory of open quantum systems proposed in Ref. 55 the correlation function $\langle \mathcal{A}_{\mu\nu} \rho_{\mu\nu} \rangle$ (with $\mu \neq \nu$, no summation over μ and ν) can be written in the form

$$\begin{aligned} \langle \mathcal{A}_{\mu\nu}(t) \rho_{\mu\nu}(t) \rangle = & \langle \mathcal{A}_{\mu\nu}^{(0)}(t) \rangle \langle \rho_{\mu\nu}(t) \rangle + \int dt_1 \langle \mathcal{A}_{\mu\nu}^{(0)}(t), \mathcal{A}_{\mu'\nu'}^{(0)+}(t_1) \rangle \\ & \times \langle i[\rho_{\mu\nu}(t), \rho_{\nu'\mu'}(t_1)]_- \rangle \theta(t - t_1) \\ & + \int dt_1 \langle i[\mathcal{A}_{\mu\nu}^{(0)}(t), \mathcal{A}_{\mu'\nu'}^{(0)+}(t_1)]_- \rangle \\ & \times \langle \rho_{\nu'\mu'}(t_1) \rho_{\mu\nu}(t) \rangle \theta(t - t_1). \end{aligned} \quad (\text{A18})$$

Here $\mathcal{A}_{\mu\nu}^{(0)}(t)$ is a variable of the free environment (with no

coupling to the electron-proton system), and $\theta(t-t_1)$ is the Heaviside unit step function. We introduce the following notations for a cumulant function of two operators $\mathcal{A}(t)$ and $\mathcal{B}(t)$:

$$\langle \mathcal{A}(t), \mathcal{B}(t') \rangle = \langle \mathcal{A}(t) \mathcal{B}(t') \rangle - \langle \mathcal{A}(t) \rangle \langle \mathcal{B}(t') \rangle,$$

and for a commutator:

$$[\mathcal{A}(t), \mathcal{B}(t')]_- = \mathcal{A}(t) \mathcal{B}(t') - \mathcal{B}(t') \mathcal{A}(t).$$

In Eq. (A18) we take into account the backaction of the bath in a contrast to the approach of Ref. 29 where this backaction is not included into consideration. Due to significant decoherence effects, off-diagonal elements of the density matrix, $\langle \rho_{\mu\nu}(t) \rangle$, disappear very fast. Because of this, the first term in the right hand side of Eq. (A18) can be neglected despite the nonzero value of the average unperturbed operator $\langle \mathcal{A}_{\mu\nu}^{(0)}(t) \rangle$. The times t and t_1 involved in the integrands of Eq. (A18) are separated by the correlation time τ_c of the correlators, which are similar to the function $\langle Q_{AL}(t), Q_{AL}(t_1) \rangle$. The time scale τ_c is determined by the reorganization energy λ_{AL} and temperature T : $\tau_c \sim \hbar / \sqrt{\lambda_{AL} T}$ [see Refs. 54 and 56 and Eq. (A23) below]. We assume that transitions between the active sites have a negligible effect on the time evolution of the operator $\rho_{\mu\nu}$ between the times t and t_1 , which are separated by the correlation time τ_c . Thus, the correlation functions and commutators of the operators $\rho_{\mu\nu}(t)$ and $\rho_{\nu\mu}(t_1)$ can be calculated using free-evolving functions,

$$\rho_{\mu\nu}(t) = \rho_{\mu\nu}(t_1) \exp\{i\omega_{\mu\nu}(t - t_1)\},$$

where $\omega_{\mu\nu} = E_\mu - E_\nu$. For correlator (A18) we obtain the formula

$$\begin{aligned} \langle \mathcal{A}_{\mu\nu}(t) \rho_{\mu\nu}(t) \rangle = & i \int dt_1 e^{i\omega_{\mu\nu}(t-t_1)} \theta(t - t_1) \{ \langle \mathcal{A}_{\mu\nu}^{(0)}(t), \mathcal{A}_{\mu\nu}^{(0)+}(t_1) \rangle \\ & \times \langle \rho_\mu(t) \rangle - \langle \mathcal{A}_{\mu\nu}^{(0)+}(t_1), \mathcal{A}_{\mu\nu}^{(0)}(t) \rangle \langle \rho_\nu(t) \rangle \}. \end{aligned} \quad (\text{A19})$$

With Eq. (A16), we can express the cumulant $\langle \mathcal{A}_{\mu\nu}^{(0)} \times(t), \mathcal{A}_{\mu\nu}^{(0)+}(t_1) \rangle$ in terms of cumulant functions of the unperturbed bath operators $Q_{AL}^{(0)}, \dots, Q_{ZD}^{(0)}$,

$$\begin{aligned} \langle \mathcal{A}_{\mu\nu}^{(0)}(t), \mathcal{A}_{\mu\nu}^{(0)+}(t_1) \rangle = & |(a_A^+ a_L)_{\mu\nu}|^2 \langle Q_{AL}^{(0)}(t), Q_{AL}^{(0)+}(t_1) \rangle \\ & + |(a_L^+ a_R)_{\mu\nu}|^2 \langle Q_{LR}^{(0)}(t), Q_{LR}^{(0)+}(t_1) \rangle \\ & + |(a_R^+ a_B)_{\mu\nu}|^2 \langle Q_{RB}^{(0)}(t), Q_{RB}^{(0)+}(t_1) \rangle \\ & + |(b_D^+ b_X)_{\mu\nu}|^2 \langle Q_{DX}^{(0)}(t), Q_{DX}^{(0)+}(t_1) \rangle \\ & + |(b_X^+ b_C)_{\mu\nu}|^2 \langle Q_{XC}^{(0)}(t), Q_{XC}^{(0)+}(t_1) \rangle \\ & + |(b_D^+ b_Z)_{\mu\nu}|^2 \langle Q_{DZ}^{(0)}(t), Q_{DZ}^{(0)+}(t_1) \rangle. \end{aligned} \quad (\text{A20})$$

The correlation function $\langle \mathcal{A}_{\mu\nu}^{(0)+}(t_1), \mathcal{A}_{\mu\nu}^{(0)}(t) \rangle$ has a similar form, with cumulants $\langle Q_{AL}^{(0)}(t), Q_{AL}^{(0)+}(t_1) \rangle, \dots$, being replaced by $\langle Q_{AL}^{(0)+}(t_1), Q_{AL}^{(0)}(t) \rangle, \dots$. Using definition (A10) of the bath operators we can calculate their correlation functions. In particular,

$$\begin{aligned}\langle Q_{AL}^{(0)}(t), Q_{AL}^{(0)+}(t_1) \rangle &= |\Delta_{AL}|^2 \exp\{-i\mathcal{W}_{AL}^{(1)}(t-t_1)\} \\ &\quad \times \exp\{-\mathcal{W}_{AL}^{(2)}(t-t_1)\}, \\ \langle Q_{AL}^{(0)+}(t_1), Q_{AL}^{(0)}(t) \rangle &= |\Delta_{AL}|^2 \exp\{i\mathcal{W}_{AL}^{(1)}(t-t_1)\} \\ &\quad \times \exp\{-\mathcal{W}_{AL}^{(2)}(t-t_1)\},\end{aligned}\quad (\text{A21})$$

where

$$\mathcal{W}_{AL}^{(1)}(\tau) = \sum_j \frac{m_j \omega_j}{2\hbar} (x_{jA} - x_{jL})^2 \sin \omega_j \tau, \quad (\text{A22})$$

$$\mathcal{W}_{AL}^{(2)}(\tau) = \sum_j \frac{m_j \omega_j}{2\hbar} \coth\left(\frac{\hbar \omega_j}{2T}\right) (x_{jA} - x_{jL})^2 (1 - \cos \omega_j \tau),$$

and T is the temperature of the environment ($k_B=1$). These expressions can be simplified in the high-temperature limit when the thermal fluctuations are much faster ($\omega_j \tau \ll 1$) than the environment modes coupled to the charge transfer:⁵⁴ $\mathcal{W}_{AL}^{(1)}(\tau) = \lambda_{AL} \tau$, $\mathcal{W}_{AL}^{(2)}(\tau) = \lambda_{AL} T \tau^2$, and correspondingly,

$$\begin{aligned}\langle Q_{AL}^{(0)}(t), Q_{AL}^{(0)+}(t_1) \rangle &= |\Delta_{AL}|^2 e^{-i\lambda_{AL}(t-t_1)} e^{-\lambda_{AL}T(t-t_1)^2}, \\ \langle Q_{AL}^{(0)+}(t_1), Q_{AL}^{(0)}(t) \rangle &= |\Delta_{AL}|^2 e^{i\lambda_{AL}(t-t_1)} e^{-\lambda_{AL}T(t-t_1)^2}.\end{aligned}\quad (\text{A23})$$

We introduce here the reorganization energy,

$$\lambda_{AL} = \sum_j \frac{m_j \omega_j^2 (x_{jA} - x_{jL})^2}{2}, \quad (\text{A24})$$

corresponding to the electron transition from site A to site L . Similar parameters can also be introduced for other electron transitions: From L to R , from R to B , as well as for proton transitions between sites D and X , X and C , and between D and the catalytic site Z .

After a sequential substitution of Eqs. (A23), (A20), and (A19) into the averaged Eq. (A17), we obtain the contribution of the intersite transfers into the master equation

$$\langle \dot{\rho}_\mu \rangle = \langle -i[\rho_\mu, H_{tr}]_- \rangle + \sum_\nu \kappa_{\mu\nu} \langle \rho_\nu \rangle - \sum_\nu \kappa_{\nu\mu} \langle \rho_\mu \rangle, \quad (\text{A25})$$

where the combined rate $\kappa_{\mu\nu}$ contains contributions of all possible electron and proton transitions,

$$\begin{aligned}\kappa_{\mu\nu} &= (\kappa_{AL})_{\mu\nu} + (\kappa_{LR})_{\mu\nu} + (\kappa_{RB})_{\mu\nu} + (\kappa_{DX})_{\mu\nu} \\ &\quad + (\kappa_{XC})_{\mu\nu} + (\kappa_{DZ})_{\mu\nu}.\end{aligned}\quad (\text{A26})$$

The rates corresponding to the specific electron transfers, $(\kappa_{AL})_{\mu\nu}$, $(\kappa_{LR})_{\mu\nu}$, $(\kappa_{RB})_{\mu\nu}$, and the rates related to the proton transfers, $(\kappa_{DX})_{\mu\nu}$, $(\kappa_{XC})_{\mu\nu}$, $(\kappa_{DZ})_{\mu\nu}$, are all determined by the Marcus equations^{54,56} with coefficients given by the appropriate transition matrices. In particular,

$$\begin{aligned}(\kappa_{AL})_{\mu\nu} &= |\Delta_{AL}|^2 \sqrt{\frac{\pi}{\lambda_{AL} T}} (|(a_A^\dagger a_L)_{\mu\nu}|^2 + |(a_A^\dagger a_L)_{\nu\mu}|^2) \\ &\quad \times \exp\left[-\frac{(E_\mu - E_\nu + \lambda_{AL})^2}{4\lambda_{AL} T}\right],\end{aligned}\quad (\text{A27})$$

$$\begin{aligned}(\kappa_{DZ})_{\mu\nu} &= |\Delta_{DZ}|^2 \sqrt{\frac{\pi}{\lambda_{DZ} T}} (|(b_D^\dagger b_Z)_{\mu\nu}|^2 + |(b_D^\dagger b_Z)_{\nu\mu}|^2) \\ &\quad \times \exp\left[-\frac{(E_\mu - E_\nu + \lambda_{DZ})^2}{4\lambda_{DZ} T}\right].\end{aligned}\quad (\text{A28})$$

It should be noted that the ratio between the transposed rate coefficients is equal to the Boltzmann factor, as

$$\frac{(\kappa_{AL})_{\mu\nu}}{(\kappa_{AL})_{\nu\mu}} = \exp\left(-\frac{E_\mu - E_\nu}{T}\right),$$

which results in the Boltzmann distribution for the equilibrium density matrix of the system.

The contribution, $\langle -i[\rho_\mu, H_{tr}]_- \rangle$, of proton transitions between site D and N -side of the membrane and between the exit site C and the P -side of the membrane to the master Eq. (A25) can be calculated with the methods of quantum transport theory.^{29,46,47} The coupling to the proton reservoirs is described by the relaxation matrix,

$$\begin{aligned}\gamma_{\mu\nu} &= \Gamma_N \{ |b_{D;\mu\nu}|^2 [1 - F_N(\omega_{\nu\mu})] + |b_{D;\nu\mu}|^2 F_N(\omega_{\mu\nu}) \} \\ &\quad + \Gamma_P \{ |b_{C;\mu\nu}|^2 [1 - F_P(\omega_{\nu\mu})] + |b_{C;\nu\mu}|^2 F_P(\omega_{\mu\nu}) \},\end{aligned}\quad (\text{A29})$$

where the energy-independent coefficient,

$$\Gamma_\sigma = 2\pi \sum_q |T_{q\sigma}|^2 \delta(\omega - \varepsilon_{q\sigma}), \quad (\text{A30})$$

determines the rate of a proton delivery to the D -site ($\sigma=N$) or the rate of a proton removal from the C -site ($\sigma=P$). We assume here that protons on the σ -side of the membrane are described by the Fermi distribution,

$$F_\sigma(\varepsilon_{q\sigma}) = \left[\exp\left(\frac{\varepsilon_{q\sigma} - \mu_\sigma}{T}\right) + 1 \right]^{-1}, \quad (\text{A31})$$

characterized by a chemical potential μ_σ .

As a result, we obtain the system of master equations for the probabilities $\langle \rho_\mu \rangle$, as

$$\langle \dot{\rho}_\mu \rangle = \sum_\nu (\kappa_{\mu\nu} + \gamma_{\mu\nu}) \langle \rho_\nu \rangle - \sum_\nu (\kappa_{\nu\mu} + \gamma_{\nu\mu}) \langle \rho_\mu \rangle, \quad (\text{A32})$$

where the intersite rates $\kappa_{\mu\nu}$ are determined by Eqs. (A26) and (A27), and the relaxation matrix, $\gamma_{\mu\nu}$, is given by Eq. (A29).

6. Algebraic solution of the master equations

Determination of the time-dependent solution of the master Eq. (A32) can be reduced to a purely algebraic problem. To accomplish this, we rewrite the Eq. (A32) in the form

$$\langle \dot{\rho}_\mu \rangle = - \sum_\nu \Lambda_{\mu\nu} \langle \rho_\nu \rangle, \quad (\text{A33})$$

with a total relaxation matrix $\Lambda_{\mu\nu}$, where $\Lambda_{\mu\nu} = -(\kappa_{\mu\nu} + \gamma_{\mu\nu})$ at $\mu \neq \nu$, and $\Lambda_{\mu\mu} = \sum_\nu (\kappa_{\nu\mu} + \gamma_{\nu\mu})$. The vector $\langle \rho \rangle$ with the elements $\langle \rho_\mu \rangle$ ($\mu=1, \dots, 64$) can be represented as a sum of the steady-state part, ρ^0 , and the time-dependent deviation $\tilde{\rho}(t)$, as $\langle \rho \rangle = \rho^0 + \tilde{\rho}$. Both the total probability vector $\langle \rho \rangle$ and

its steady-state value satisfy the normalization condition: $\sum_{\mu} \langle \rho_{\mu} \rangle = \sum_{\mu} \rho_{\mu}^0 = 1$. The steady-state distribution can be found from the matrix equation $\Lambda \rho^0 = 0$, and for a time-dependent part $\tilde{\rho}$ we have a rate equation in the form $(d/dt)\tilde{\rho} = -\Lambda \tilde{\rho}$. Using the unitary operator, \mathcal{U} , the matrix Λ can be transformed to the diagonal form $\Lambda' = \mathcal{U}^{-1} \Lambda \mathcal{U}$ with γ'_{μ} as the diagonal elements. This transformation should be accompanied by the transformation of the vector $\tilde{\rho}$ as $\tilde{\rho} = \mathcal{U} \rho'$. Then, the vector $\rho'(t)$ obeys the diagonal equation with a simple solution for its μ -component: $\rho'_{\mu}(t) = e^{-\gamma'_{\mu} t} \rho'_{\mu}(0)$. Correspondingly, the time evolution of the probability vector $\rho(t)$ from its initial value $\rho(0)$ is described by the formula

$$\langle \rho(t) \rangle = \rho^0 - \mathcal{S}(t) \rho^0 + \mathcal{S}(t) \rho(0), \quad (\text{A34})$$

where $\mathcal{S}(t) = \mathcal{U} \mathcal{Z}(t) \mathcal{U}^{-1}$, and $\mathcal{Z}(t)$ is the diagonal matrix with the elements $\mathcal{Z}_{\mu\nu}(t) = \delta_{\mu\nu} e^{-\gamma'_{\mu} t}$. It should be noted that $\mathcal{S}(0) = \hat{1}$ and $\mathcal{S}(\infty) = 0$, where $\hat{1}$ is the 64×64 unit matrix.

7. Proton current

The time-dependent populations, $\langle n_{\alpha} \rangle$ and $\langle N_{\beta} \rangle$, of all redox and protonable sites in the model are expressed in terms of the evolving probability distribution $\langle \rho(t) \rangle$. We recall that the index α labels the redox sites $\alpha = A$ (Cu_A), L (heme a), R (heme a_3), and B (Cu_B). The index β labels the protonable sites $\beta = D, X, C$, and Z . Finally, the index σ labels the two sides of the membrane $\sigma = N, P$. With the density matrix probability distributions, $\langle \rho_{\mu}(t) \rangle$, over the states, $|\mu\rangle$, of the system we can also find the proton flows from the N - and P -sides of the membrane into the system, $I_{\sigma} = (d/dt)\langle N_{\sigma} \rangle$, where $\langle N_{\sigma} \rangle$ is the total number of protons on the σ -side of the membrane: $\langle N_{\sigma} \rangle = \sum_q \langle N_{q\sigma} \rangle$. Using techniques developed in quantum transport theory,^{29,46} we obtain the formulas for the proton currents I_N and I_P ,

$$I_N = \Gamma_N \sum_{\mu\nu} |b_{D;\mu\nu}|^2 \{ [1 - F_N(\omega_{\nu\mu})] \langle \rho_{\nu} \rangle - F_N(\omega_{\nu\mu}) \langle \rho_{\mu} \rangle \}, \quad (\text{A35})$$

$$I_P = \Gamma_P \sum_{\mu\nu} |b_{C;\mu\nu}|^2 \{ [1 - F_P(\omega_{\nu\mu})] \langle \rho_{\nu} \rangle - F_P(\omega_{\nu\mu}) \langle \rho_{\mu} \rangle \}.$$

Note that these currents depend on the time-dependent probability distribution $\langle \rho(t) \rangle$ and, accordingly, they also vary with time. The total number of protons, $\langle N_{\sigma} \rangle$, transferred to the σ -side of the membrane ($\sigma = P, N$) is calculated as the integral of the corresponding current,

$$\langle N_{\sigma}(t) \rangle = \int_0^t dt_1 I_{\sigma}(t_1). \quad (\text{A36})$$

8. Proton-motive force

The proton-motive force across the membrane can be defined as a difference of electrochemical potentials μ_P and μ_N involved in the Fermi distribution (A31) of the proton reservoirs: $\Delta\mu = \mu_P - \mu_N$. This gradient includes the transmembrane concentration difference (ΔpH) and the transmembrane voltage V ,

$$\Delta\mu = V - 2.3(RT/F) \times \Delta pH. \quad (\text{A37})$$

Here R and F are the gas and Faraday constant, respectively, and T is the temperature (in degrees Kelvin, $k_B = 1$).^{1,2} Both energy parameters, $\Delta\mu$ and V , are measured in meV. At the standard conditions ($T = 298$ K), the concentration gradient contributes about 60 meV per ΔpH -unit. This results in the transmembrane voltage $V \sim 150$ meV, provided that the total proton-motive force, $\Delta\mu$, is about 210 meV, and $\Delta pH = -1$.¹

The transmembrane voltage, $V > 0$, elevates the energies of protonable sites adjacent to the P -side and lowers the energies of the proton sites located near the N -side.²² The electron sites are simultaneously experiencing the opposite effect, for the same V . As a result the electron energy levels, ε_{α} , and the proton energies, ε_{β} , involved in the Hamiltonian H_0 (A1) are shifted from their initial values, $\varepsilon_{\alpha}^{(0)}$ and $\varepsilon_{\beta}^{(0)}$,

$$\begin{aligned} \varepsilon_{\alpha} &= \varepsilon_{\alpha}^{(0)} - V(x_{\alpha}/W), \\ \varepsilon_{\beta} &= \varepsilon_{\beta}^{(0)} + V(x_{\beta}/W), \end{aligned} \quad (\text{A38})$$

where W is the membrane width. The positions of the redox and protonable sites, x_{α} and x_{β} , are counted here from the middle of the membrane with the x -axis directed toward the P -side: $x_A \sim W/2$, $x_L \sim x_R \sim x_B \sim W/6$, $x_D \sim 0.1W$, $x_X \sim 0.3W$, $x_C = 0.5W$, $x_Z = W/6$.^{8,10}

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